

# **Time-domain ab initio analysis of facet-dependent** carrier dynamics in Cuprous oxide



Minjae Kwen<sup>1</sup>, Yoosang Son<sup>1</sup>, and Hyungjun Kim<sup>1\*</sup>

<sup>1</sup>Department of Chemistry, Korea Advanced Institute of Science and Technology

## Introduction

Cuprous oxide ( $Cu_2O$ ) is a promising photoactive material for photoelectrodes used in water splitting and CO2 reduction.<sup>[1,2]</sup> A through characterization of charge carrier mobility and recombination dynamics in Cu<sub>2</sub>O is important to crucial for understanding its photoelectrochemical (PEC) behavior.

Performance of the Cu<sub>2</sub>O photoelectrode was studied by *Pan et al.*<sup>[3]</sup> Carrier dynamics of single crystalline Cu<sub>2</sub>O films were examined using ultrafast transient reflection spectroscopy.



## Results

Electronic Structure of Cu<sub>2</sub>O



Figures from Fig 2a, 2d, 2e of *Nature* 628.8009 (2024): 765-770.<sup>[3]</sup>

#### **Exciton lifetime and carrier diffusion lengths follows:**



## **Calculational Methods**

Surface hopping method with classical path approximation (CPA) was used for nonadiabatic dynamics simulation.



Electronic band structures are achieved by using VASP and VASPKIT<sup>[6]</sup>

### Surface Hopping Simulation



#### [Lifetime of exciton (ps)]

Bulk	(100)	(110)	(111)
Cu <sub>2</sub> O	facet	facet	facet
148	0.730	16.1	22.4

### Classical Path Approximation (CPA)<sup>[4]</sup>

The CPA is the assumption that the nuclear dynamics is unaffected by the electronic dynamics. With CPA, the nuclear trajectory is predetermined during the surface hopping simulation.



### Decoherence Induced Surface Hopping (DISH)<sup>[5]</sup>

(1) Time propagation of waveftn.  $\rightarrow$  Coefficient calculation governed by TDSE Using non-adiabatic coupling (NAC),  $H_{ij} \equiv -i\hbar \left\langle \Phi_i \left| \frac{\partial}{\partial t} \right| \Phi_j \right\rangle$ 

(2) Decoherence time  $(\tau_m)$  calculation, from the autocorrelation ftn. of energy gap

$$\frac{1}{\tau_{\alpha}}(t) = \sum_{i \neq \alpha}^{N} |c_{i}(t)|^{2} r_{\alpha i}, \quad r_{\alpha i} \text{ is the time constant of the decay of } J_{ij}(t) = \langle \Psi_{i}(t) | \Psi_{j}(t) \rangle$$
$$J_{ij}(t) \approx \langle J_{ij}(\mathbf{R}(t)) \rangle_{T} \approx \exp\left[-\frac{1}{\hbar^{2}} \int_{0}^{t} d\tau_{2} \int_{0}^{\tau_{2}} d\tau_{1} \langle \delta U(t) \cdot \delta U(t_{0}) \rangle_{T}\right], \quad \delta U(t) = \Delta E_{ij}(\mathbf{R}(t)) - \langle \Delta E_{ij}(\mathbf{R}(t)) \rangle_{T}$$



Surface hopping simulation performed using Quantum Espresso and PYXAID<sup>[4]</sup>

Decomposed partial charge density







### Conclusion

Using semiclassical nonadiabatic dynamics simulations, we elucidated the quantum mechanisms behind the slower carrier recombination observed at (111) facet. We also explicated this phenomenon with high band separation of VBM and CBM at (111) slab system.

#### References

[1] Paracchino, Adriana, et al. "Highly active oxide photocathode for photoelectrochemical water reduction". Nature materials 10.6 (2011), pp. 456-461. [2] Li, Christina W., and Matthew W. Kanan. "CO2 reduction at low overpotential on Cu electrodes resulting from the reduction of thick Cu2O films". Journal of the American Chemical Society 134.17 (2012), pp. 7231-7234.

[3] Pan, Linfeng, et al. "High carrier mobility along the [111] orientation in Cu2O photoelectrodes". Nature 628.8009 (2024), pp. 765-770. [4] Akimov, Alexey V., and Oleg V. Prezhdo. "The PYXAID program for non-adiabatic molecular dynamics in condensed matter systems". Journal of chemical theory and computation 9.11 (2013), pp. 4959-4972.

[5] Jaeger, Heather M., et al. "Decoherence-induced surface hopping." The Journal of chemical physics 137.22 (2012). [6] Wang, V., et al. "A user-friendly interface facilitating high-throughput computing and analysis using VASP code., 2021, 267."